

Fast Methods for Simulation of Biomolecule Electrostatics

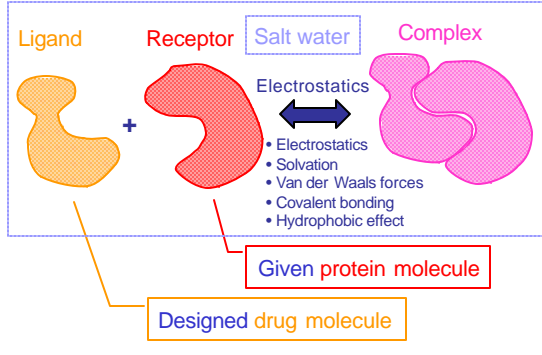
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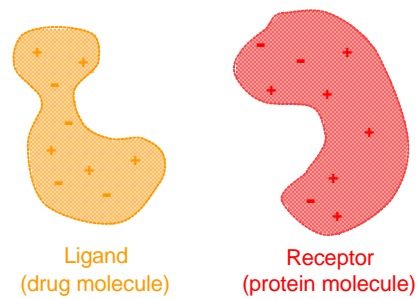
Outline

- > Problem statement
- > Finite-difference approach and problems
- > Integral equation method and advantages
- > Fast solver implementation
- > Computational results
- > Conclusion and Future work

The Problem of Drug Design

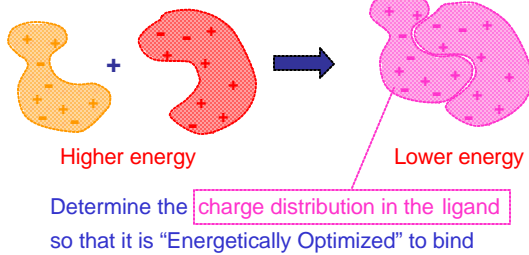


Electrostatics View

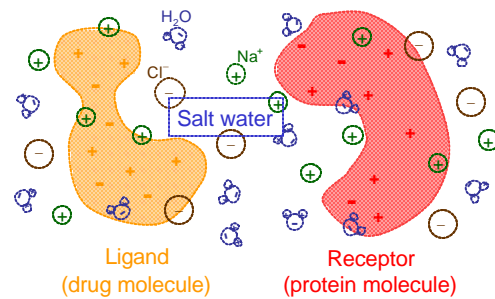


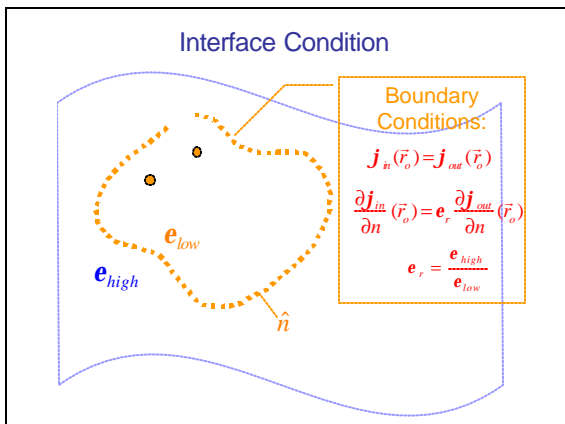
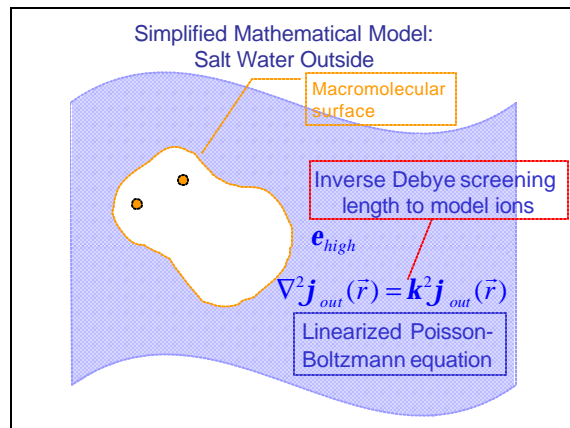
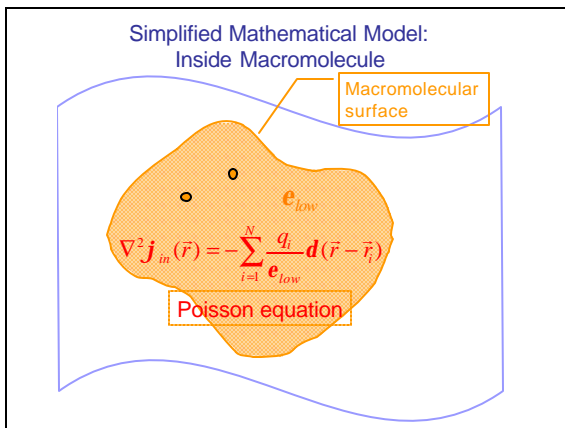
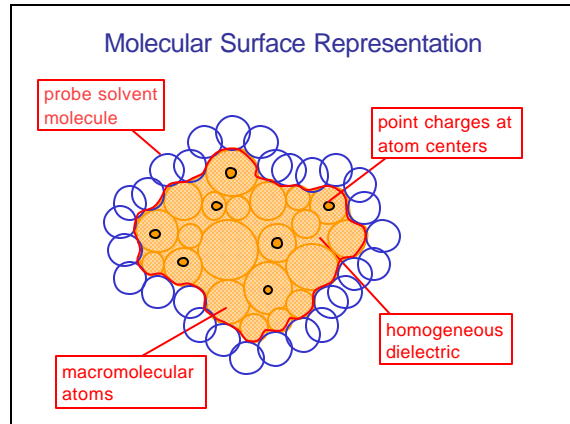
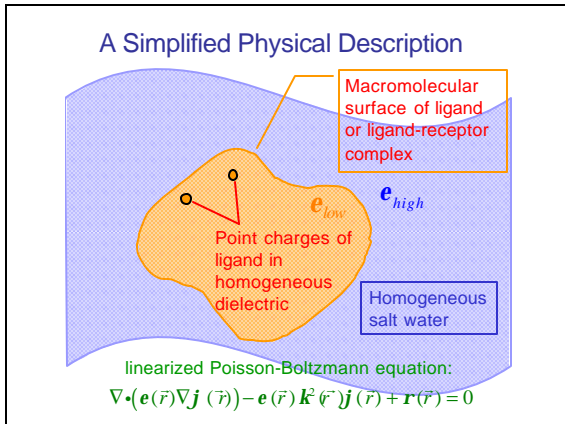
Minimize Electrostatic Binding Energy

$$E_{binding} = |E_{desolvation}^{ligand}| + |E_{desolvation}^{receptor}| - |E_{interaction}|$$



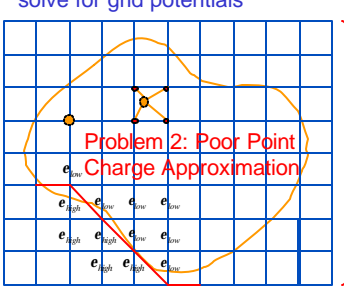
An Electrostatic Analysis Problem





- ### Why Use this Simplified Model?
- Atomistic Level Simulation is too expensive
 - Salt ions and water molecules treated individually
 - Continuum Model Matches Well with Experimental Data

Standard Finite-Difference Method
 set up boundary conditions and solve for grid potentials

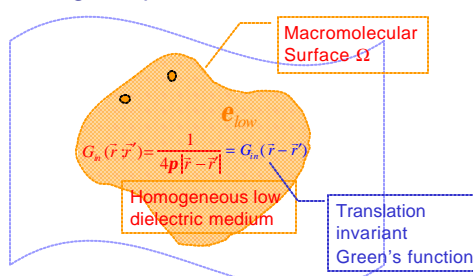


Problem 1: Inaccurate Molecular Surface

Problem 2: Poor Point Charge Approximation

Problem 3: Inexact Boundary Conditions

Integral equation: Interior Problem



Macromolecular Surface Ω

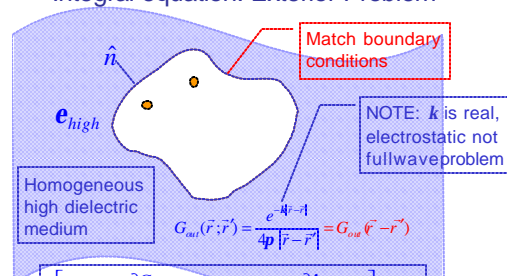
Homogeneous low dielectric medium

Translation invariant Green's function

$$G_a(\vec{r}, \vec{r}') = \frac{1}{4\pi|\vec{r} - \vec{r}'|} = G_a(\vec{r} - \vec{r}')$$

$$\int_{\Omega} \left[\mathbf{j}_{in}(\vec{r}') \frac{\partial G_{in}(\vec{r}; \vec{r}')}{\partial n} - G_{in}(\vec{r}; \vec{r}') \frac{\partial \mathbf{j}_{in}(\vec{r}')}{\partial n} \right] d\vec{r}' = \sum_{k=1}^N \frac{q_k}{\epsilon_{low}} G_{in}(\vec{r}; \vec{r}_k)$$

Integral equation: Exterior Problem



Match boundary conditions

NOTE: k is real, electrostatic not fullwave problem

Homogeneous high dielectric medium

$$G_{out}(\vec{r}; \vec{r}') = \frac{e^{-k|\vec{r} - \vec{r}'|}}{4\pi|\vec{r} - \vec{r}'|} = G_{out}(\vec{r} - \vec{r}')$$

$$\int_{\Omega} \left[-\mathbf{j}_{out}(\vec{r}') \frac{\partial G_{out}(\vec{r}; \vec{r}')}{\partial n} + G_{out}(\vec{r}; \vec{r}') \frac{\partial \mathbf{j}_{out}(\vec{r}')}{\partial n} \right] d\vec{r}' = 0$$

$\mathbf{j}_{in}(\vec{r}') = \frac{1}{\epsilon_r} \frac{\partial \mathbf{j}_{in}(\vec{r}')}{\partial n}$

Advantages For Integral Equation Formulation

- Directly discretize surfaces
- Point charges treated exactly
- Handles infinite exterior

$$\int_{\Omega} \left[\mathbf{j}_{in}(\vec{r}') \frac{\partial G_{in}(\vec{r}; \vec{r}')}{\partial n} - G_{in}(\vec{r}; \vec{r}') \frac{\partial \mathbf{j}_{in}(\vec{r}')}{\partial n} \right] d\vec{r}' = \sum_{k=1}^N \frac{q_k}{\epsilon_{low}} G_{in}(\vec{r}; \vec{r}_k)$$

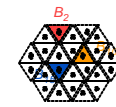
$$\int_{\Omega} \left[-\mathbf{j}_{in}(\vec{r}') \frac{\partial G_{out}(\vec{r}; \vec{r}')}{\partial n} + G_{out}(\vec{r}; \vec{r}') \frac{1}{\epsilon_r} \frac{\partial \mathbf{j}_{in}(\vec{r}')}{\partial n} \right] d\vec{r}' = 0$$

Standard piecewise constant collocation discretization method

$$\mathbf{j}_{in}(\vec{r}) \approx \sum_j a_j B_j(\vec{r})$$

$$\frac{\partial \mathbf{j}_{in}(\vec{r})}{\partial n} \approx \sum_j b_j B_j(\vec{r})$$

$\vec{r} \in \Omega$



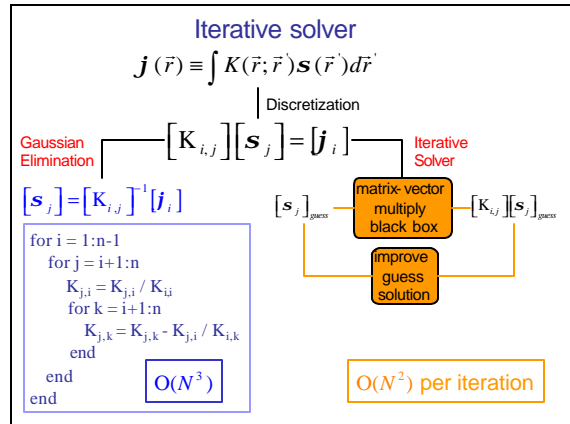
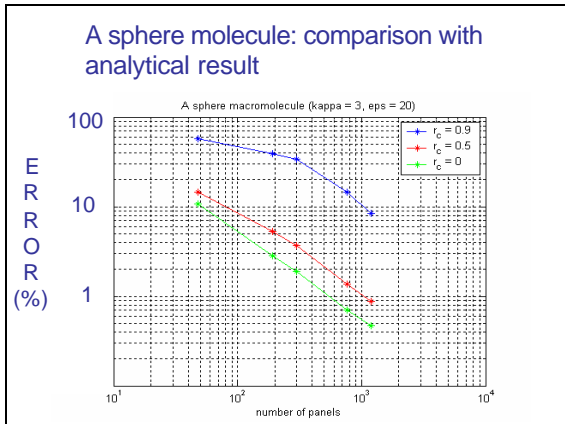
- Piecewise constant basis functions
- Collocation points at panel centroids

Matrix Equation

$$\begin{bmatrix} D^{in} & S^h \\ D^{out} & S^{out} \end{bmatrix} \begin{bmatrix} a_j \\ b_j \end{bmatrix} = \begin{bmatrix} \sum_{k=1}^N \frac{q_k}{4\pi|\vec{r}_i - \vec{r}_k|} \\ 0 \end{bmatrix}$$

$$D_{ij}^{in} = \int_{panel_j} \frac{\partial}{\partial n'} \left(\frac{1}{4\pi|\vec{r}_i - \vec{r}'|} \right) d\vec{r}' \quad S_j^{in} = - \int_{panel_j} \frac{1}{4\pi|\vec{r}_i - \vec{r}'|} d\vec{r}'$$

$$D_{ij}^{out} = - \int_{panel_j} \frac{\partial}{\partial n'} \left(\frac{e^{-k|\vec{r}_i - \vec{r}'|}}{4\pi|\vec{r}_i - \vec{r}'|} \right) d\vec{r}' \quad S_j^{out} = \frac{1}{\epsilon_r} \int_{panel_j} \frac{e^{-k|\vec{r}_i - \vec{r}'|}}{4\pi|\vec{r}_i - \vec{r}'|} d\vec{r}'$$



Use Fast Integral Equation Solver

$O(N \log N)$ Matrix-vector multiply

- Multiple Green's functions
- Translation Invariant kernel

$$G_{in}(\vec{r}, \vec{r}') = \frac{1}{4\pi |\vec{r} - \vec{r}'|}$$

$$G_{out}(\vec{r}, \vec{r}') = \frac{e^{-k|\vec{r} - \vec{r}'|}}{4\pi |\vec{r} - \vec{r}'|}$$

Pre-corrected FFT algorithm

$$[\mathbf{s}_j]_{\text{guess}} \rightarrow \text{matrix-vector multiply black box} \rightarrow [\mathbf{K}_{i,j}] [\mathbf{s}_j]_{\text{guess}}$$

charge distribution \rightarrow pre-corrected FFT \rightarrow potential due to space-invariant kernel

- projection of panel charges onto grid charges
- grid potentials due to grid charges are computed by FFT
- potentials on panel centroids are interpolated from grid potentials
- direct interaction and correction among near neighbors

Picture courtesy of J. Phillips

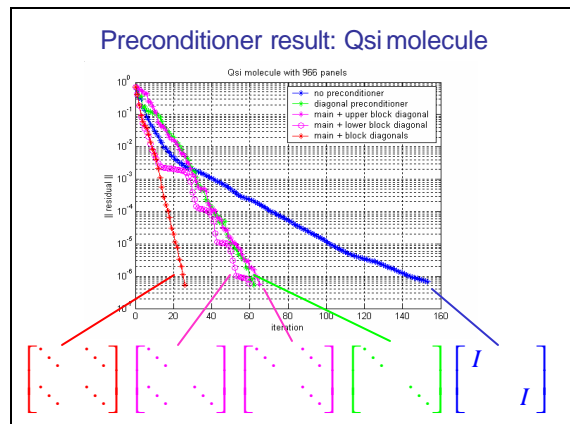
Preconditioner on Two Examples

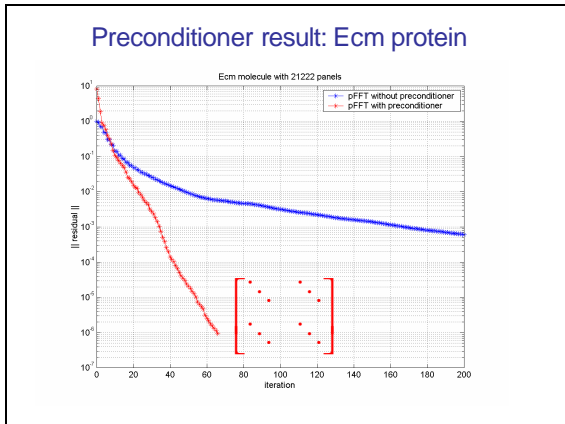
Need to find a good preconditioner $[P] \approx [\mathbf{K}_{i,j}]$

And solve $[P]^{-1} [\mathbf{K}_{i,j}] [\mathbf{s}_j] = [P]^{-1} [\mathbf{j}_i]$

hopefully better conditioned than $[\mathbf{K}_{i,j}]$

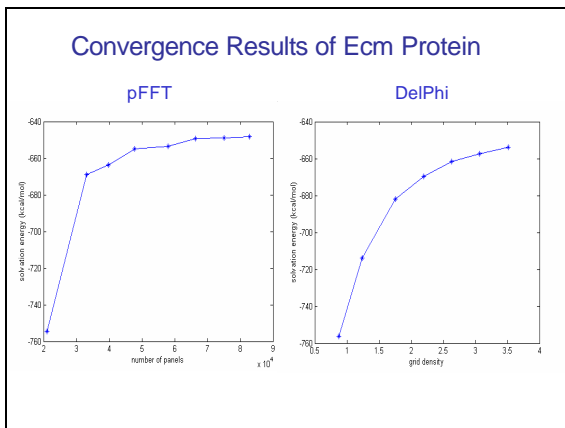
Qsi molecule Ecm protein





Accuracy comparison with DelPhi

	# of dielectric panels	# of salt panels	$E_{\text{solvation}}$ (kcal/mol)	
			pFFT	DelPhi
Water	17204	9330	-3.14	-3.17
TSA	34114	5842	-34.62	-34.75
ECM	82868	18596	-646.42	-653.88



Binding energy calculation of a protein-peptide complex

	Energy calculated (kcal/mol)			
	$R_{\text{desolvation}}$	$L_{\text{desolvation}}$	$(R \rightarrow L)_{\text{interaction}}$	$(L \rightarrow R)_{\text{interaction}}$
pFFT	14.52	24.47	130.80	130.91
DelPhi	14.51	24.47	131.03	131.03

Conclusions and Future work

- > Carefully selected **Integral Formulation** results in **Fast Solver** for Biomolecule Electrostatics
- > Working on coupling to **charge optimization problem** in drug design
- > Extending formulation to include more complicated geometry (**inner cavities** in macromolecule)
- > Fine tuning existing **pre-corrected FFT code**